

One-Parameter Irreducible Representation of $\text{spl}(2, 1)$ Superalgebra

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Abstract One-parameter irreducible representation of the $\text{spl}(2, 1)$ superalgebra is studied on subspace and quotient spaces of the universal enveloping algebra of Heisenberg-Weyl superalgebra. The parameter α may be related to the interaction parameter U in one exactly solvable model for correlated electrons.

Keywords $\text{spl}(2, 1)$ superalgebra · Irreducible representation · Exactly solvable model

1 Introduction

A series of models of correlated electrons on a lattice and exactly solvable in one dimension and supersymmetric, such as Hubbard and extended Hubbard models and t - J model, EKS model, BGLZ model [1], has been extensively studied due to their promising role in theoretical condensed-matter physics and possibly in high- T_c superconductivity. Those models contain one symmetry-preserving free real parameter which is the Hubbard interaction parameter U . One-parameter irreducible representations of Lie superalgebra have played an important role in constructing supersymmetrical models. The supersymmetrical algebra of BGLZ model for correlated electrons on the unrestricted 4^L -dimensional electronic Hilbert space $\bigotimes_{n=1}^L \mathbb{C}^4$ is superalgebra $\text{gl}(2|1)$. Some irreducible representations of Lie superalgebras $\text{spl}(2, 1)$ and $\text{gl}(2|1)$ have been given by Chen [2–4]. One-parameter indecomposable representation of the $\text{spl}(2, 1)$ superalgebra has been studied [5]. In the present paper we shall be concerned with the $\text{spl}(2, 1)$ superalgebra. The purpose of the present paper is to study new one-parameter irreducible representation of the $\text{spl}(2, 1)$ superalgebra on subspace and quotient spaces of the universal enveloping algebra of Heisenberg-Weyl superalgebra in terms of indecomposable representation of this superalgebra.

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2 One-Parameter Irreducible Representation of the $\text{spl}(2, 1)$

In accordance with Chen [2] the generators of the $\text{spl}(2, 1)$ superalgebra read as follows:

$$\{Q_3, Q_+, Q_-, B \in \text{spl}(2, 1)_{\bar{0}} \mid V_+, V_-, W_+, W_- \in \text{spl}(2, 1)_{\bar{1}}\} \quad (1)$$

and satisfy the following commutation and anticommutation relations:

$$\begin{aligned} [Q_3, Q_{\pm}] &= \pm Q_{\pm}, & [Q_+, Q_-] &= 2Q_3, & [B, Q_{\pm}] &= [B, Q_3] = 0 \\ [Q_3, V_{\pm}] &= \pm \frac{1}{2}V_{\pm}, & [Q_3, W_{\pm}] &= \pm \frac{1}{2}W_{\pm}, & [B, V_{\pm}] &= \frac{1}{2}V_{\pm} \\ [B, W_{\pm}] &= -\frac{1}{2}W_{\pm}, & [Q_{\pm}, V_{\mp}] &= V_{\pm}, & [Q_{\pm}, W_{\mp}] &= W_{\pm}, & [Q_{\pm}, V_{\pm}] &= 0 \\ [Q_{\pm}, W_{\pm}] &= 0, & \{V_{\pm}, V_{\pm}\} &= \{V_{\pm}, V_{\mp}\} = \{W_{\pm}, W_{\pm}\} = \{W_{\pm}, W_{\mp}\} &= 0, \\ \{V_{\pm}, W_{\pm}\} &= \pm Q_{\pm}, & \{V_{\pm}, W_{\mp}\} &= -Q_3 \pm B. \end{aligned} \quad (2)$$

In Ref. [5] we have obtained one-parameter indecomposable representation of the $\text{spl}(2, 1)$ superalgebra.

The generalized Fock space is defined as a quotient space of V

$$Y = (V/J) : \{\phi(k, \alpha_1, \alpha_2) = \phi(k, 0, \alpha_1, 0, \alpha_2, 0) \bmod J \mid k \in \mathbb{Z}^+, \alpha_1, \alpha_2 = 0, 1\} \quad (3)$$

where J is the left ideal generated by the element $b - \lambda, a_1 - \eta_1$ and $a_2 - \eta_2$, λ is a complex number and η_1 and η_2 are generators of the Grassmann algebra \tilde{G} . On this space, a new representation form can be obtained as follows

$$\begin{aligned} L(Q_3)\phi(k, \alpha_1, \alpha_2) &= \left(-\frac{1}{2}n + k + \frac{1}{2}\alpha_1 + \frac{1}{2}\alpha_2\right)\phi(k, \alpha_1, \alpha_2), \\ L(Q_+)\phi(k, \alpha_1, \alpha_2) &= (n - k - \alpha_1 - \alpha_2)\phi(k + 1, \alpha_1, \alpha_2), \\ L(Q_-)\phi(k, \alpha_1, \alpha_2) &= k\phi(k - 1, \alpha_1, \alpha_2), \\ L(B)\phi(k, \alpha_1, \alpha_2) &= \left[\left(\frac{1}{2} + \alpha\right)n - \frac{1}{2}\alpha_1 - \frac{1}{2}\alpha_2\right]\phi(k, \alpha_1, \alpha_2), \\ L(V_+)\phi(k, \alpha_1, \alpha_2) &= \alpha_1\sqrt{\alpha}\phi(k + 1, \alpha_1 - 1, \alpha_2) \\ &\quad + (-1)^{\alpha_1}(1 - \alpha_2)(n - k - \alpha_1)\sqrt{1 + \alpha}\phi(k, \alpha_1, \alpha_2 + 1), \\ L(V_-)\phi(k, \alpha_1, \alpha_2) &= \alpha_1\sqrt{\alpha}\phi(k, \alpha_1 - 1, \alpha_2) \\ &\quad - (-1)^{\alpha_1}(1 - \alpha_2)\sqrt{1 + \alpha}k\phi(k - 1, \alpha_1, \alpha_2 + 1), \\ L(W_+)\phi(k, \alpha_1, \alpha_2) &= (-1)^{\alpha_1}\alpha_2\sqrt{1 + \alpha}\phi(k + 1, \alpha_1, \alpha_2 - 1) \\ &\quad + (-n + k + \alpha_2)(1 - \alpha_1)\sqrt{\alpha}\phi(k, \alpha_1 + 1, \alpha_2), \\ L(W_-)\phi(k, \alpha_1, \alpha_2) &= (1 - \alpha_1)\sqrt{\alpha}k\phi(k - 1, \alpha_1 + 1, \alpha_2) \\ &\quad + (-1)^{\alpha_1}\alpha_2\sqrt{1 + \alpha}\phi(k, \alpha_1, \alpha_2 - 1). \end{aligned} \quad (4)$$

We can easily see that the representation (4) is an infinite-dimensional irreducible representation when $n \notin Z^+$. Obviously, the invariant subspace exists when $n \in Z^+$,

$$\begin{aligned} Y(n) : & \{\phi(k, \alpha_1, \alpha_2) \in Y \mid k + \alpha_1 + \alpha_2 \leq n, k \in Z^+, \alpha_1, \alpha_2 = 0, 1\}, \\ & \dim Y(n) = 4n, \end{aligned} \quad (5)$$

and there is no invariant complementary subspace. Thus, the representation (4) is indecomposable. Restricting the representation given by (4) to the invariant subspace $Y(n)$, we can obtain a finite-dimensional irreducible representation of the $\text{spl}(2, 1)$.

For the sake of simplicity, we redefine the basis of $Y(n)$ as

$$\begin{aligned} |j, m, \alpha_1, \alpha_2\rangle = & \sqrt{\frac{(j-m)!(2j-\alpha_1-\alpha_2)!}{(j+m)!(j-m-\alpha_1)!(j-m-\alpha_2)!}} \\ & \times \phi(j+m, \alpha_1, \alpha_2), \end{aligned} \quad (6)$$

where $j = \frac{1}{2}n = 0, \frac{1}{2}, 1, \frac{3}{2}, 2, \dots$,

$$\begin{aligned} m = -j, -j+1, \dots, j, & \quad \text{when } \alpha_1 = 0, \alpha_2 = 0, \\ m = -j, -j+1, \dots, j-1, & \quad \text{when } \alpha_1 = 0, \alpha_2 = 1, \\ m = -j, -j+1, \dots, j-1, & \quad \text{when } \alpha_1 = 1, \alpha_2 = 0, \\ m = -j, -j+1, \dots, j-2, & \quad \text{when } \alpha_1 = 1, \alpha_2 = 1. \end{aligned} \quad (7)$$

The action of the generators of the $\text{spl}(2, 1)$ on the new basis vector is straightforwardly obtained with the help of (4) and (6). One finds

$$\begin{aligned} Q_3|j, m, \alpha_1, \alpha_2\rangle &= \left(m + \frac{1}{2}\alpha_1 + \frac{1}{2}\alpha_2\right)|j, m, \alpha_1, \alpha_2\rangle, \\ Q_+|j, m, \alpha_1, \alpha_2\rangle &= (j-m-\alpha_1-\alpha_2)\sqrt{\frac{(j+m+1)(j-m)}{(j-m-\alpha_1)(j-m-\alpha_2)}}|j, m+1, \alpha_1, \alpha_2\rangle, \\ Q_-|j, m, \alpha_1, \alpha_2\rangle &= \sqrt{\frac{(j+m)(j-m+1-\alpha_1)(j-m+1-\alpha_2)}{j-m+1}}|j, m-1, \alpha_1, \alpha_2\rangle, \\ B|j, m, \alpha_1, \alpha_2\rangle &= \left[(1+2\alpha)j - \frac{1}{2}\alpha_1 - \frac{1}{2}\alpha_2\right]|j, m, \alpha_1, \alpha_2\rangle, \\ V_+|j, m, \alpha_1, \alpha_2\rangle &= (-1)^{\alpha_1}(1-\alpha_2)\sqrt{1+\alpha}(j-m-\alpha_1)\frac{1}{\sqrt{j-m-\alpha_2}}|j, m, \alpha_1, \alpha_2+1\rangle \\ &+ \alpha_1\sqrt{\alpha}\sqrt{\frac{(j+m+1)(j-m)}{j-m-\alpha_2}}|j, m+1, \alpha_1-1, \alpha_2\rangle, \\ V_-|j, m, \alpha_1, \alpha_2\rangle &= \alpha_1\sqrt{\alpha}\sqrt{j-m+1-\alpha_1}|j, m, \alpha_1-1, \alpha_2\rangle \\ &- (-1)^{\alpha_1}(1-\alpha_2)\sqrt{1+\alpha}\sqrt{\frac{j+m}{(j-m+1)(j-m-\alpha_2)}} \\ &\times |j, m-1, \alpha_1, \alpha_2+1\rangle \end{aligned} \quad (8)$$

$$\begin{aligned}
W_+|j, m, \alpha_1, \alpha_2\rangle &= -(1 - \alpha_1)\sqrt{\alpha}(j - m - \alpha_2) \frac{1}{\sqrt{j - m - \alpha_1}} |j, m, \alpha_1 + 1, \alpha_2\rangle \\
&\quad + (-1)^{\alpha_1} \alpha_2 \sqrt{1 + \alpha} \sqrt{\frac{(j + m + 1)(j - m)}{j - m - \alpha_1}} |j, m + 1, \alpha_1, \alpha_2 - 1\rangle \\
W_-|j, m, \alpha_1, \alpha_2\rangle &= (1 - \alpha_1)\sqrt{\alpha} \sqrt{\frac{(j + m)(j - m + 1 - \alpha_2)}{j - m + 1}} |j, m - 1, \alpha_1 + 1, \alpha_2\rangle \\
&\quad + (-1)^{\alpha_1} \alpha_2 \sqrt{1 + \alpha} \sqrt{j - m + 1 - \alpha_2} |j, m, \alpha_1, \alpha_2 - 1\rangle,
\end{aligned}$$

where we restrict $|j, j + 1, \alpha_1, \alpha_2\rangle = |j, -j - 1, \alpha_1, \alpha_2\rangle = 0$.

3 Conclusion

To illustrate the irreducibility of the $\text{spl}(2, 1)$ representation, we have a simple discussion. In the first place, the representation space $Y(2j)$ set up by all $|j, m, \alpha_1, \alpha_2\rangle$ marked with j is invariant under the action of the $\text{spl}(2, 1)$ generators. In the next place, there is no true subspace in the $Y(2j)$. It is clear from (8) that this representation is a $8j$ -dimensional irreducible representation.

We have obtained one-parameter irreducible representation. All the finite-dimensional one-parameter irreducible representation of the $\text{spl}(2, 1)$ have been given on the subspace of the generalized Fock space. In terms of the conclusion it may be of use for further researches on one-parameter coherent state of the $\text{spl}(2, 1)$ superalgebra and for determining new supersymmetrical quantum model corresponding to the $\text{spl}(2, 1)$.

Appendix

Here we derive the new representation form (4) in terms of one-parameter indecomposable representation of the $\text{spl}(2, 1)$ superalgebra. In following space

$$V: \{\phi(k, l, \alpha_1, \beta_1, \alpha_2, \beta_2) | k, l \in Z^+, \alpha_1, \beta_1, \alpha_2, \beta_2 = 0, 1\} \quad (\text{A.1})$$

its explicit form as follows:

$$\begin{aligned}
L(Q_3)\phi(k, l, \alpha_1, \beta_1, \alpha_2, \beta_2) &= \left(-\frac{1}{2}n + k + \frac{1}{2}\alpha_1 + \frac{1}{2}\alpha_2\right) \phi(k, l, \alpha_1, \beta_1, \alpha_2, \beta_2) + \phi(k + 1, l + 1, \alpha_1, \beta_1, \alpha_2, \beta_2) \\
&\quad + \frac{1}{2}(-1)^{\alpha_1}(1 - \alpha_1)\phi(k, l, \alpha_1 + 1, \beta_1 + 1, \alpha_2, \beta_2) \\
&\quad + \frac{1}{2}(-1)^{\alpha_2}(1 - \alpha_2)\phi(k, l, \alpha_1, \beta_1, \alpha_2 + 1, \beta_2 + 1),
\end{aligned}$$

$$\begin{aligned}
L(Q_+)\phi(k, l, \alpha_1, \beta_1, \alpha_2, \beta_2) &= (n - k - \alpha_1 - \alpha_2)\phi(k + 1, l, \alpha_1, \beta_1, \alpha_2, \beta_2) - \phi(k + 2, l + 1, \alpha_1, \beta_1, \alpha_2, \beta_2) \\
&\quad - (-1)^{\alpha_1}(1 - \alpha_1)\phi(k + 1, l, \alpha_1 + 1, \beta_1 + 1, \alpha_2, \beta_2)
\end{aligned}$$

$$\begin{aligned}
& -(-1)^{\alpha_2}(1-\alpha_2)\phi(k+1,l,\alpha_1,\beta_1,\alpha_2+1,\beta_2+1), \\
L(Q_-)\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) &= \phi(k,l+1,\alpha_1,\beta_1,\alpha_2,\beta_2)+k\phi(k-1,l,\alpha_1,\beta_1,\alpha_2,\beta_2), \\
L(B)\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) &= \left[\left(\frac{1}{2} + \alpha \right) n - \frac{1}{2}\alpha_1 - \frac{1}{2}\alpha_2 \right] \phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) \\
& + \frac{1}{2}(-1)^{\alpha_1}(1-\alpha_1)\phi(k,l,\alpha_1+1,\beta_1+1,\alpha_2,\beta_2) \\
& + \frac{1}{2}(-1)^{\alpha_2}(1-\alpha_2)\phi(k,l,\alpha_1,\beta_1,\alpha_2+1,\beta_2+1), \\
L(V_+)\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) &= (-1)^{\alpha_1+\beta_1}(1-\alpha_2)(n-k-\alpha_1)\sqrt{1+\alpha}\phi(k,l,\alpha_1,\beta_1,\alpha_2+1,\beta_2) \\
& + (-1)^{\alpha_1}\sqrt{\alpha}\phi(k+1,l,\alpha_1,\beta_1+1,\alpha_2,\beta_2) \\
& + \alpha_1\sqrt{\alpha}\phi(k+1,l,\alpha_1-1,\beta_1,\alpha_2,\beta_2) \\
& - (-1)^{\alpha_1+\beta_1}(1-\alpha_2)\sqrt{1+\alpha}\phi(k+1,l+1,\alpha_1,\beta_1,\alpha_2+1,\beta_2) \\
& - (-1)^{\beta_1}(1-\alpha_1)(1-\alpha_2)\sqrt{1+\alpha}\phi(k,l,\alpha_1+1,\beta_1+1,\alpha_2+1,\beta_2), \\
L(V_-)\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) &= (-1)^{\alpha_1}\sqrt{\alpha}\phi(k,l,\alpha_1,\beta_1+1,\alpha_2,\beta_2) + \alpha_1\sqrt{\alpha}\phi(k,l,\alpha_1-1,\beta_1,\alpha_2,\beta_2) \\
& - (-1)^{\alpha_1+\beta_1}(1-\alpha_2)\sqrt{1+\alpha}\phi(k,l+1,\alpha_1,\beta_1,\alpha_2+1,\beta_2) \\
& - (-1)^{\alpha_1+\beta_1}(1-\alpha_2)\sqrt{1+\alpha}k\phi(k-1,l,\alpha_1,\beta_1,\alpha_2+1,\beta_2), \\
L(W_+)\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) &= (-n+k+\alpha_2)(1-\alpha_1)\sqrt{\alpha}\phi(k,l,\alpha_1+1,\beta_1,\alpha_2,\beta_2) \\
& + (-1)^{\alpha_1+\beta_1+\alpha_2}\sqrt{1+\alpha}\phi(k+1,l,\alpha_1,\beta_1,\alpha_2,\beta_2+1) \\
& + (-1)^{\alpha_1+\beta_1}\alpha_2\sqrt{1+\alpha}\phi(k+1,l,\alpha_1,\beta_1,\alpha_2-1,\beta_2) \\
& + (1-\alpha_1)\sqrt{\alpha}\phi(k+1,l+1,\alpha_1+1,\beta_1,\alpha_2,\beta_2) \\
& + (-1)^{\alpha_2}(1-\alpha_1)(1-\alpha_2)\sqrt{\alpha}\phi(k,l,\alpha_1+1,\beta_1,\alpha_2+1,\beta_2+1), \\
L(W_-)\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2) &= (-1)^{\alpha_1+\beta_1+\alpha_2}\sqrt{1+\alpha}\phi(k,l,\alpha_1,\beta_1,\alpha_2,\beta_2+1) \\
& + (-1)^{\alpha_1+\beta_1}\alpha_2\sqrt{1+\alpha}\phi(k,l,\alpha_1,\beta_1,\alpha_2-1,\beta_2) \\
& + (1-\alpha_1)\sqrt{\alpha}\phi(k,l+1,\alpha_1+1,\beta_1,\alpha_2,\beta_2) \\
& + (1-\alpha_1)\sqrt{\alpha}k\phi(k-1,l,\alpha_1+1,\beta_1,\alpha_2,\beta_2).
\end{aligned} \tag{A.2}$$

The generalized Fock space is defined as a quotient space of V

$$Y = (V/J) : \{\phi(k, \alpha_1, \alpha_2) = \phi(k, 0, \alpha_1, 0, \alpha_2, 0) \bmod J \mid k \in Z^+, \alpha_1, \alpha_2 = 0, 1\} \tag{A.3}$$

where J is the left ideal generated by the element $b - \lambda, \alpha_1 - \eta_1$ and $\alpha_2 - \eta_2$, λ is a complex number and η_1 and η_2 are generators of the Grassmann algebra \tilde{G} . On this space, the

representation (A.2) induces the new representation

$$\begin{aligned}
& L(Q_3)\phi(k, \alpha_1, \alpha_2) \\
&= \left(-\frac{1}{2}n + k + \frac{1}{2}\alpha_1 + \frac{1}{2}\alpha_2 \right) \phi(k, \alpha_1, \alpha_2) + \lambda\phi(k+1, \alpha_1, \alpha_2) \\
&\quad + \frac{1}{2}(1-\alpha_1)\eta_1\phi(k, \alpha_1+1, \alpha_2) + \frac{1}{2}(1-\alpha_2)\eta_2\phi(k, \alpha_1, \alpha_2+1), \\
& L(Q_+)\phi(k, \alpha_1, \alpha_2) \\
&= (n-k-\alpha_1-\alpha_2)\phi(k+1, \alpha_1, \alpha_2) \\
&\quad - \lambda\phi(k+2, \alpha_1, \alpha_2) - (1-\alpha_1)\eta_1\phi(k+1, \alpha_1+1, \alpha_2) \\
&\quad - (1-\alpha_2)\eta_2\phi(k+1, \alpha_1, \alpha_2+1), \\
& L(Q_-)\phi(k, \alpha_1, \alpha_2) = \lambda\phi(k, \alpha_1, \alpha_2) + k\phi(k-1, \alpha_1, \alpha_2), \\
& L(B)\phi(k, \alpha_1, \alpha_2) \\
&= \left[\left(\frac{1}{2} + \alpha \right) n - \frac{1}{2}\alpha_1 - \frac{1}{2}\alpha_2 \right] \phi(k, \alpha_1, \alpha_2) - \frac{1}{2}(1-\alpha_1)\eta_1\phi(k, \alpha_1+1, \alpha_2) \\
&\quad - \frac{1}{2}(1-\alpha_2)\eta_2\phi(k, \alpha_1, \alpha_2+1), \\
& L(V_+)\phi(k, \alpha_1, \alpha_2) \\
&= (-1)^{\alpha_1} \sqrt{\alpha} \eta_1\phi(k+1, \alpha_1, \alpha_2) \\
&\quad + (-1)^{\alpha_1} (1-\alpha_2) (n-k-\alpha_1) \sqrt{1+\alpha} \phi(k, \alpha_1, \alpha_2+1) \\
&\quad - (-1)^{\alpha_1} (1-\alpha_2) \sqrt{1+\alpha} \lambda\phi(k+1, \alpha_1, \alpha_2+1) \\
&\quad - (1-\alpha_1) (1-\alpha_2) \sqrt{1+\alpha} \eta_1\phi(k, \alpha_1+1, \alpha_2+1) \\
&\quad + \alpha_1 \sqrt{\alpha} \phi(k+1, \alpha_1-1, \alpha_2), \tag{A.4} \\
& L(V_-)\phi(k, \alpha_1, \alpha_2) \\
&= (-1)^{\alpha_1} \sqrt{\alpha} \eta_1\phi(k, \alpha_1, \alpha_2) - (-1)^{\alpha_1} (1-\alpha_2) \sqrt{1+\alpha} \lambda\phi(k, \alpha_1, \alpha_2+1) \\
&\quad - (-1)^{\alpha_1} (1-\alpha_2) \sqrt{1+\alpha} k\phi(k-1, \alpha_1, \alpha_2+1) \\
&\quad + \alpha_1 \sqrt{\alpha} \phi(k, \alpha_1-1, \alpha_2), \\
& L(W_+)\phi(k, \alpha_1, \alpha_2) \\
&= (-1)^{\alpha_1} \alpha_2 \sqrt{1+\alpha} \phi(k+1, \alpha_1, \alpha_2-1) \\
&\quad + (-n+k+\alpha_2)(1-\alpha_1) \sqrt{\alpha} \phi(k, \alpha_1+1, \alpha_2) \\
&\quad + (-1)^{\alpha_1+\alpha_2} \sqrt{1+\alpha} \eta_2\phi(k+1, \alpha_1, \alpha_2) \\
&\quad + (1-\alpha_1) \sqrt{\alpha} \lambda\phi(k+1, \alpha_1+1, \alpha_2) + (1-\alpha_1)(1-\alpha_2) \sqrt{\alpha} \eta_2\phi(k, \alpha_1+1, \alpha_2+1), \\
& L(W_-)\phi(k, \alpha_1, \alpha_2) \\
&= (-1)^{\alpha_1+\alpha_2} \sqrt{1+\alpha} \eta_2\phi(k, \alpha_1, \alpha_2) + (-1)^{\alpha_1} \alpha_2 \sqrt{1+\alpha} \phi(k, \alpha_1, \alpha_2-1) \\
&\quad + (1-\alpha_1) \sqrt{\alpha} \lambda\phi(k, \alpha_1+1, \alpha_2) + (1-\alpha_1) \sqrt{\alpha} k\phi(k-1, \alpha_1+1, \alpha_2).
\end{aligned}$$

The representation given by (A.4) is an infinite-dimensional irreducible representation for the cases $\lambda \neq 0, \eta_1 \neq 0$ or $\eta_2 \neq 0$. When $\lambda = \eta_1 = \eta_2 = 0$, the representation (A.4) becomes

$$\begin{aligned}
L(Q_3)\phi(k, \alpha_1, \alpha_2) &= \left(-\frac{1}{2}n + k + \frac{1}{2}\alpha_1 + \frac{1}{2}\alpha_2\right)\phi(k, \alpha_1, \alpha_2), \\
L(Q_+)\phi(k, \alpha_1, \alpha_2) &= (n - k - \alpha_1 - \alpha_2)\phi(k + 1, \alpha_1, \alpha_2), \\
L(Q_-)\phi(k, \alpha_1, \alpha_2) &= k\phi(k - 1, \alpha_1, \alpha_2), \\
L(B)\phi(k, \alpha_1, \alpha_2) &= \left[\left(\frac{1}{2} + \alpha\right)n - \frac{1}{2}\alpha_1 - \frac{1}{2}\alpha_2\right]\phi(k, \alpha_1, \alpha_2), \\
L(V_+)\phi(k, \alpha_1, \alpha_2) &= \alpha_1\sqrt{\alpha}\phi(k + 1, \alpha_1 - 1, \alpha_2) \\
&\quad + (-1)^{\alpha_1}(1 - \alpha_2)(n - k - \alpha_1)\sqrt{1 + \alpha}\phi(k, \alpha_1, \alpha_2 + 1), \\
L(V_-)\phi(k, \alpha_1, \alpha_2) &= \alpha_1\sqrt{\alpha}\phi(k, \alpha_1 - 1, \alpha_2) \\
&\quad - (-1)^{\alpha_1}(1 - \alpha_2)\sqrt{1 + \alpha}k\phi(k - 1, \alpha_1, \alpha_2 + 1), \\
L(W_+)\phi(k, \alpha_1, \alpha_2) &= (-1)^{\alpha_1}\alpha_2\sqrt{1 + \alpha}\phi(k + 1, \alpha_1, \alpha_2 - 1) \\
&\quad + (-n + k + \alpha_2)(1 - \alpha_1)\sqrt{\alpha}\phi(k, \alpha_1 + 1, \alpha_2), \\
L(W_-)\phi(k, \alpha_1, \alpha_2) &= (1 - \alpha_1)\sqrt{\alpha}k\phi(k - 1, \alpha_1 + 1, \alpha_2) + (-1)^{\alpha_1}\alpha_2\sqrt{1 + \alpha}\phi(k, \alpha_1, \alpha_2 - 1).
\end{aligned} \tag{A.5}$$

References

1. Brachen, A.J., Gould, M.D., Links, J.R., Zhang, Y.Z.: New supersymmetric and exactly solvable model of correlated electrons. *Phys. Rev. Lett.* **74**, 2768 (1995). doi:[10.1103/PhysRevLett.74.2768](https://doi.org/10.1103/PhysRevLett.74.2768)
2. Chen, Y.-Q.: Differential realizations, boson-fermion realizations of the $\text{sp}(2, 1)$ superalgebra and its representations. *J. Phys. A: Math. Gen.* **26**, 4319 (1993). doi:[10.1088/0305-4470/26/17/037](https://doi.org/10.1088/0305-4470/26/17/037)
3. Chen, Y.-Q.: One-parameter inhomogeneous differential realizations and boson-fermion realizations of the $\text{gl}(2|1)$ superalgebra. *Int. J. Theor. Phys.* **40**(6), 1113 (2001). doi:[10.1023/A:1017505619536](https://doi.org/10.1023/A:1017505619536)
4. Chen, Y.-Q.: One-parameter indecomposable and irreducible representations of $\text{gl}(2|1)$ superalgebra. *Int. J. Theor. Phys.* **40**(7), 1249 (2001). doi:[10.1023/A:1017575604884](https://doi.org/10.1023/A:1017575604884)
5. Chen, Y.-Q.: One-parameter indecomposable representation of the $\text{sp}(2, 1)$ superalgebra. *Int. J. Theor. Phys.* **45**(9), 1649 (2006). doi:[10.1007/s10773-006-9041-3](https://doi.org/10.1007/s10773-006-9041-3)